

Chiral spin liquid phase of the triangular lattice Hubbard model

Evidence from iDMRG in a mixed real- and momentum-space basis

AARON SZASZ

(With Johannes Motruk, Michael P. Zaletel, and Joel E. Moore)

UC BERKELEY

LBL

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Outline

1. Introduction
2. Calculation methods
3. Phase diagram
4. Chiral spin liquid phase
5. Implications/comparisons and summary
6. Future directions

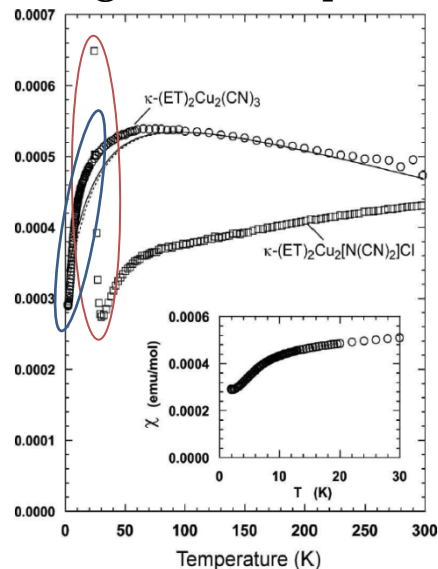
Motivation from experiments

Possible spin liquids in triangular lattice systems!

Eg: $\kappa - (ET)_2(Cu)_2(CN)_3$: approximately isotropic triangular lattice

Nonmagnetic at low T:

Magnetic susceptibility



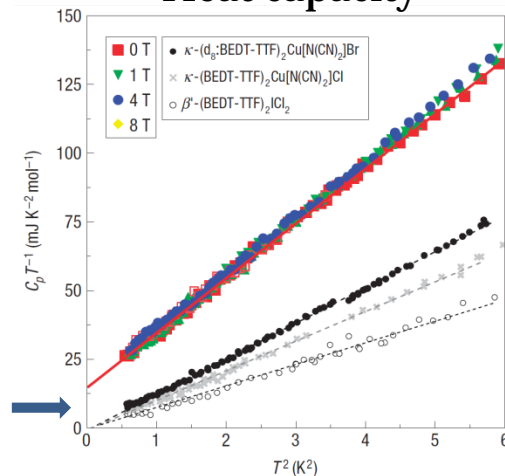
■ AFM

■ Spin liquid?

Shimizu et al., PRL 2003

Gapless?

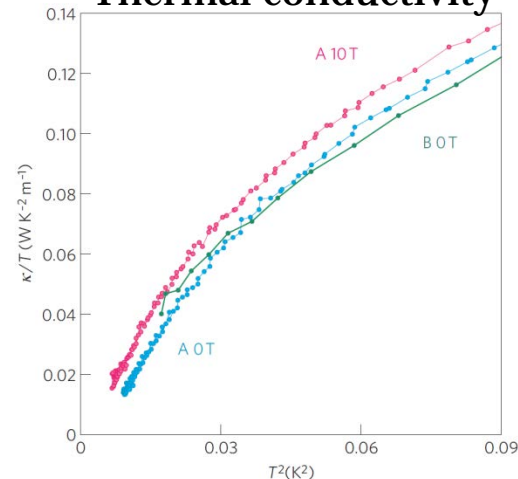
Heat capacity



Yamashita et al., Nat. Phys. 2008

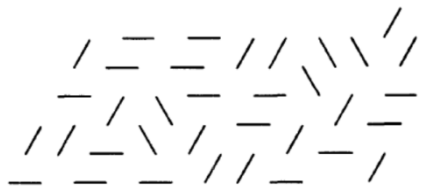
Gapped?

Thermal conductivity



Yamashita et al., Nat. Phys. 2008

Spin liquids

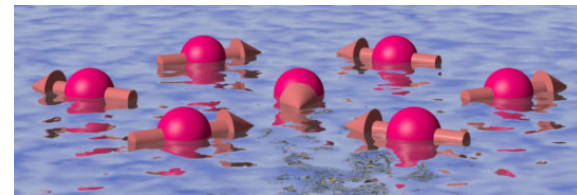


Anderson 1973

Candidate states:

Gapless states:

- U(1) spin liquid: spinon Fermi surface
- Dirac spin liquid: gapless Dirac cones at specific points in Brillouin zone
- Quadratic band-touching



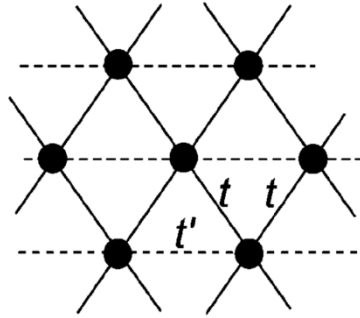
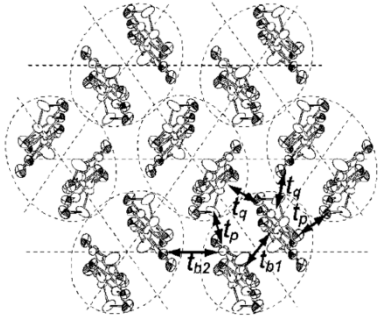
Pratt 2011

Gapped states:

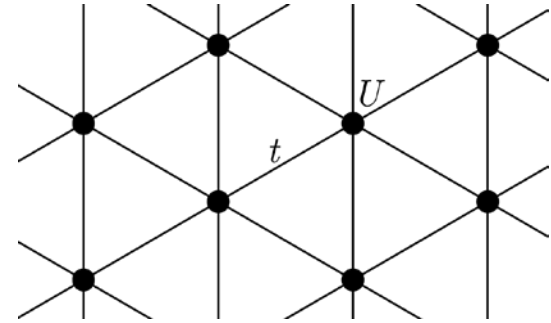
- Z_2 spin liquid: equivalent to toric code
- Chiral spin liquid:
 - Time-reversal symmetry breaking
 - Gapless edge modes
 - More later!

Models for real materials

Crystal structure of $\kappa - (ET)_2(Cu)_2(CN)_3$:



Hubbard model:



Parameter estimates:

- $t'/t = 1.06$, $U/t = 8.2$ [Shimizu et al., PRL 2003]
- $t'/t = 0.8$, $U/t \approx 12-15$ [Nakamura et al., J. Phys. Soc. Jpn. 2009]

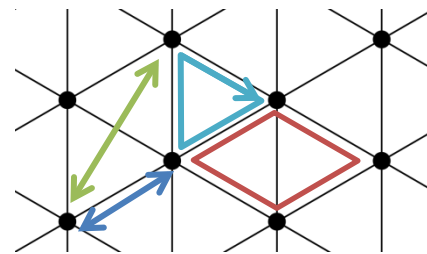
Try to simplify: t/U expansion

$1/2$ filling \rightarrow extended Heisenberg model

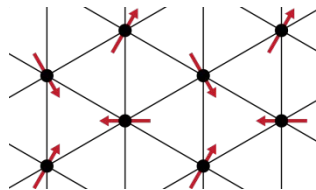
Existing results – theory

Spin models:

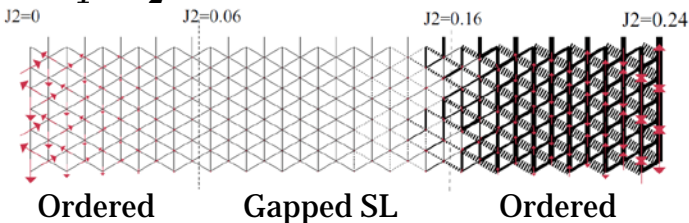
$$H = J_1 \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + J_2 \sum_{\langle\langle ij \rangle\rangle} \mathbf{S}_i \cdot \mathbf{S}_j + K \sum_{\langle ijkl \rangle} (P_{ijkl} + \text{H.c.}) + \chi \sum_{\Delta} \mathbf{S} \cdot (\mathbf{S} \times \mathbf{S})$$



Heisenberg:



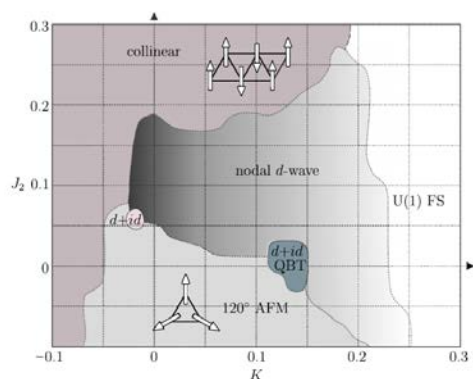
$J_1 - J_2$:



J_1, J_2, K :

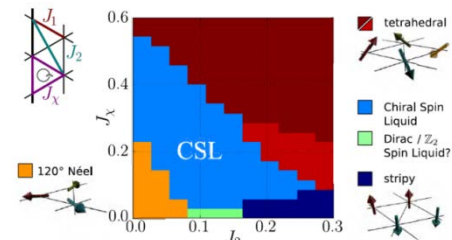
Spinon Fermi surface

- Variational: Motrunich, PRB, 2005
- Ladder DMRG: Sheng, PRB 2009

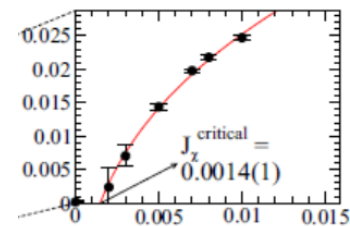


Mishmash et al, PRL 2013

J_1, J_2, χ :



Wietek & Lauchli, PRB 2017



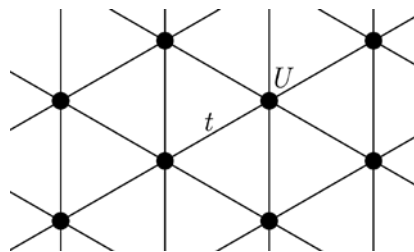
Saadatmand & McCulloch, PRB 2017

Zhu & White, PRB 2015

Existing results – theory

Hubbard model:

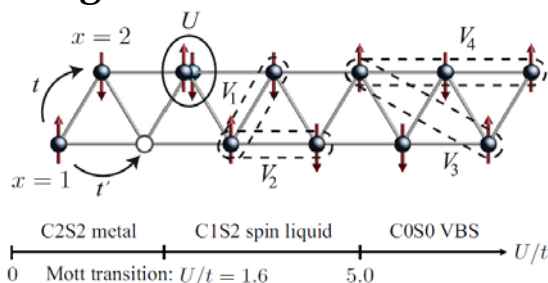
$$H = -t \sum_{\langle i,j \rangle, \sigma} \left(c_{i\sigma}^\dagger c_{j\sigma} + \text{h.c.} \right) + U \sum_i n_{i\uparrow} n_{i\downarrow}$$



Via spin models:

- Low order expansion: U(1) SL, spinon Fermi surface [Motrunich 2005]
- 12th order: SL nature unclear, maybe U(1) [Yang et al, PRL 2010]

Two leg ladder DMRG + extra terms:

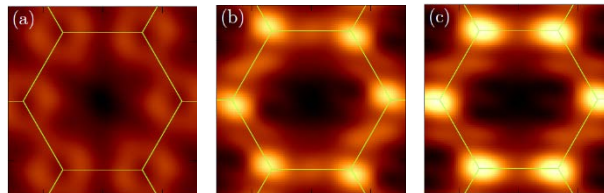
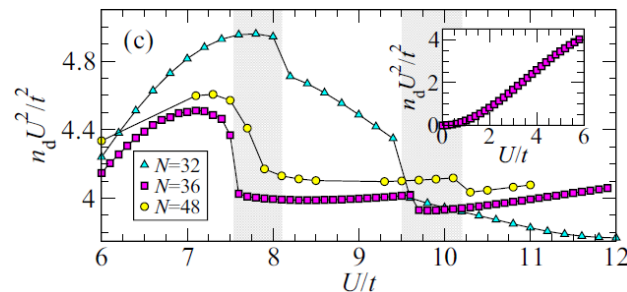


Mishmash et al, PRB 2015

DMRG on finite cylinder [RIKEN group]:

Three phases:

- Metal, SL, ordered
- First order transitions: $(\partial E / \partial U)(U/t)^2$:



Shirakawa et al, PRB 2017

SL spin correlations:

- No Fermi surface
- Nodal gapless unclear

Partial information on nature of spin liquid

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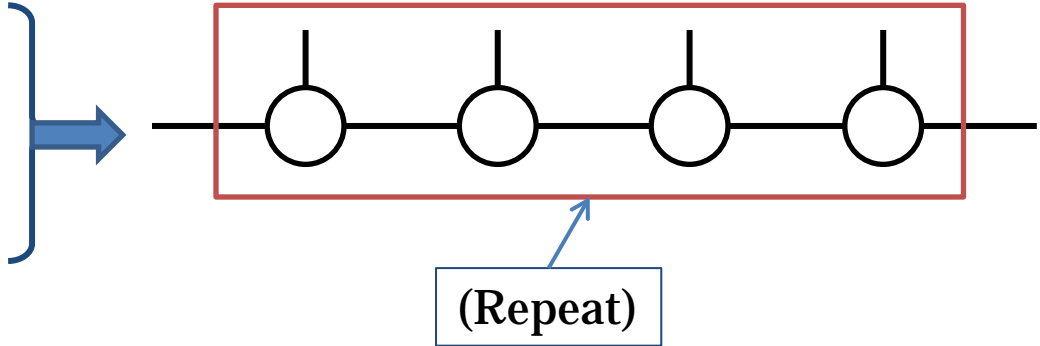
Calculation methods – iDMRG

Find ground state with the **infinite-system density matrix renormalization group (iDMRG)** method

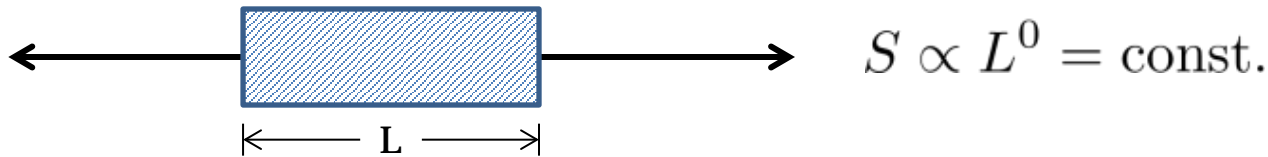
- Variational method within **Matrix Product State (MPS)** ansatz

$$|\psi\rangle = \sum_{\{\sigma_i\}} \dots A_i^{(\sigma_i)} A_{i+1}^{(\sigma_{i+1})} \dots |\dots\sigma_i\sigma_{i+1}\dots\rangle$$

$A_i \rightarrow d \times \chi \times \chi$ tensor,
d: physical dimension
 χ : MPS bond dimension

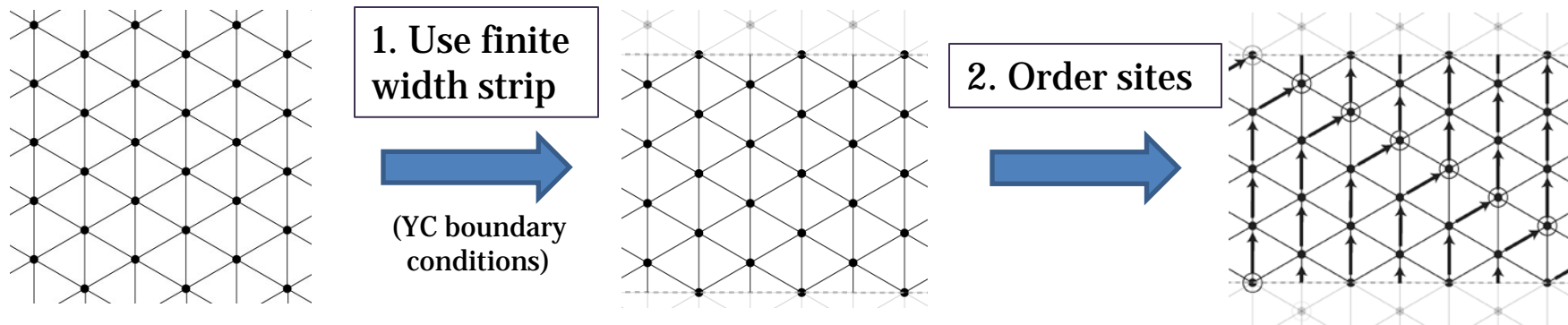


- Max entanglement $\rightarrow \log(\chi)$
 - Necessary χ scales with $\exp(S)$
 - MPS is efficient for 1D *area law* states (eg. gapped ground states!)



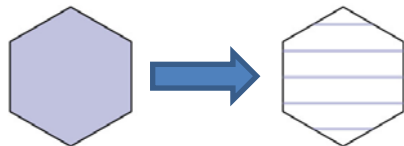
Calculation methods – cylinder DMRG

Apply DMRG to 2D system:

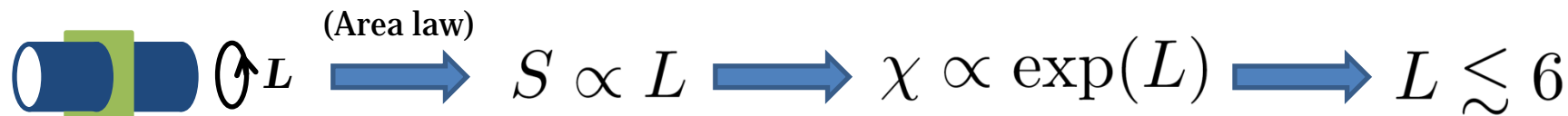


Consequences & Limitations:

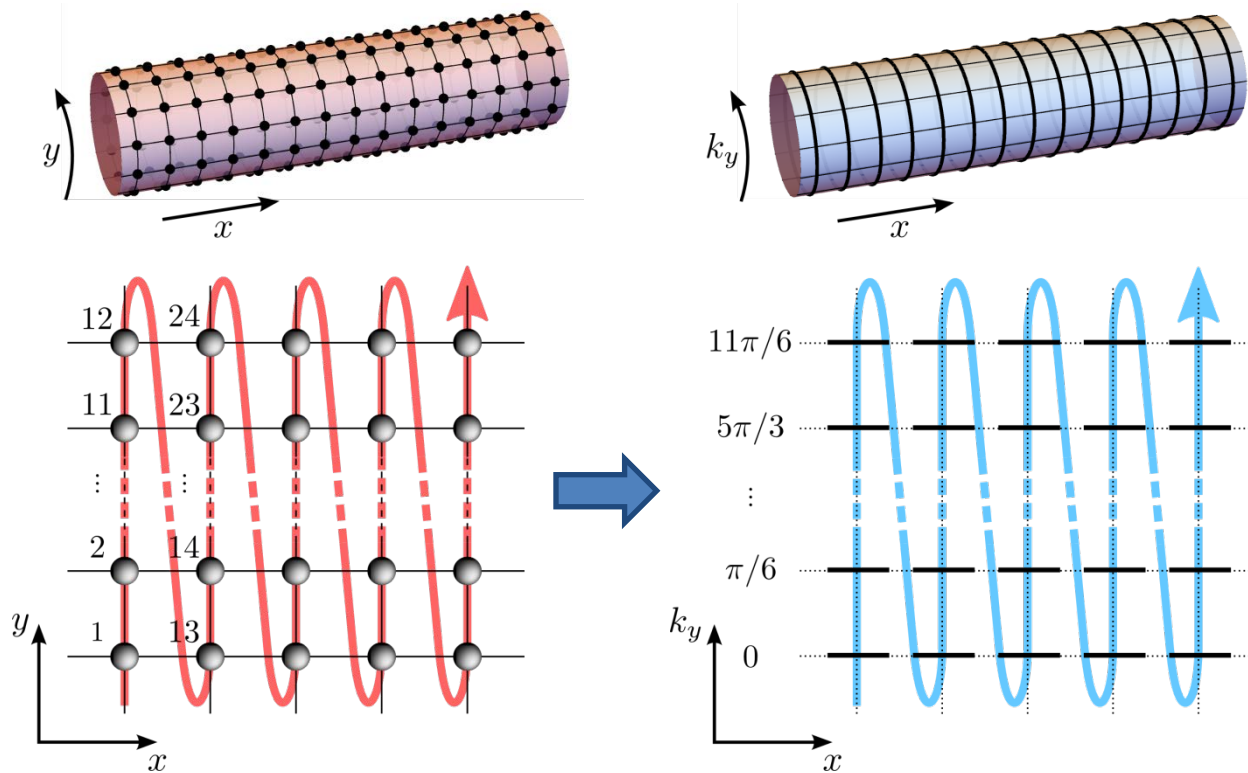
1. Discrete momenta



2. Entanglement growth

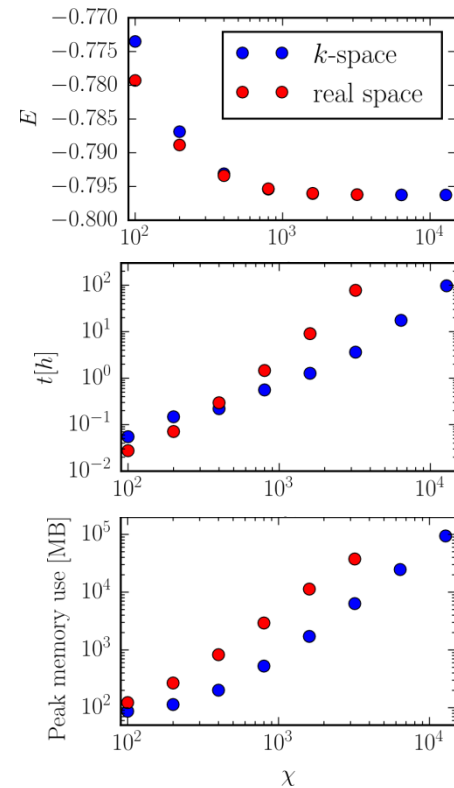


Calculation methods – mixed-space DMRG



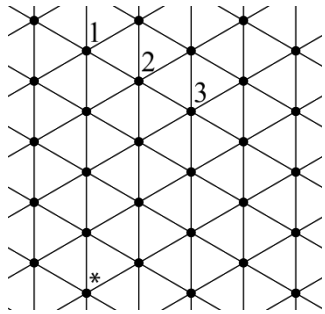
Motruk et al, PRB 2016

Benchmark: Hofstadter model

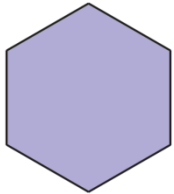


Calculation methods – mixed-space DMRG

On triangular lattice

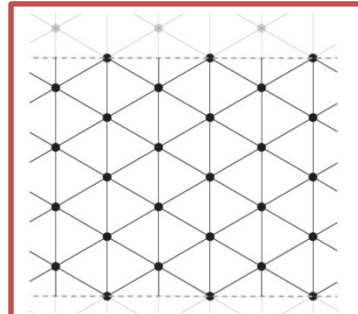


$L=4$

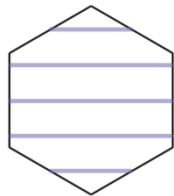


Number of k_y
quantum numbers:

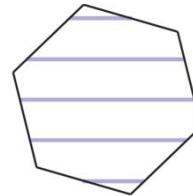
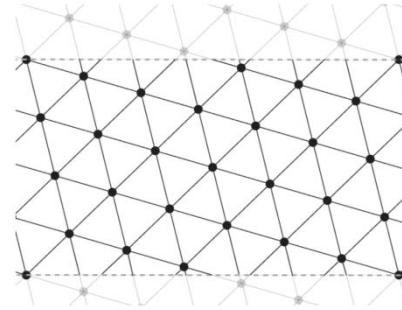
Number of k quantum numbers depends on boundary conditions:



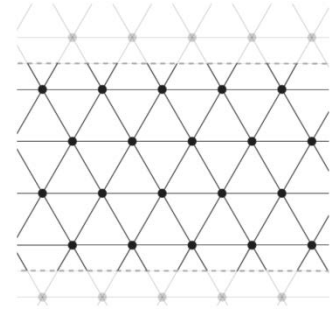
YC



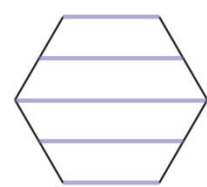
4
(L)



1



XC



2
($L/2$)

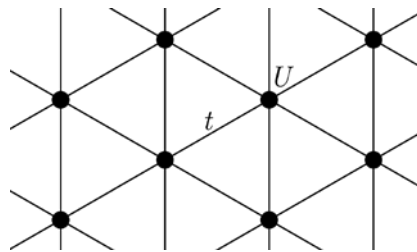
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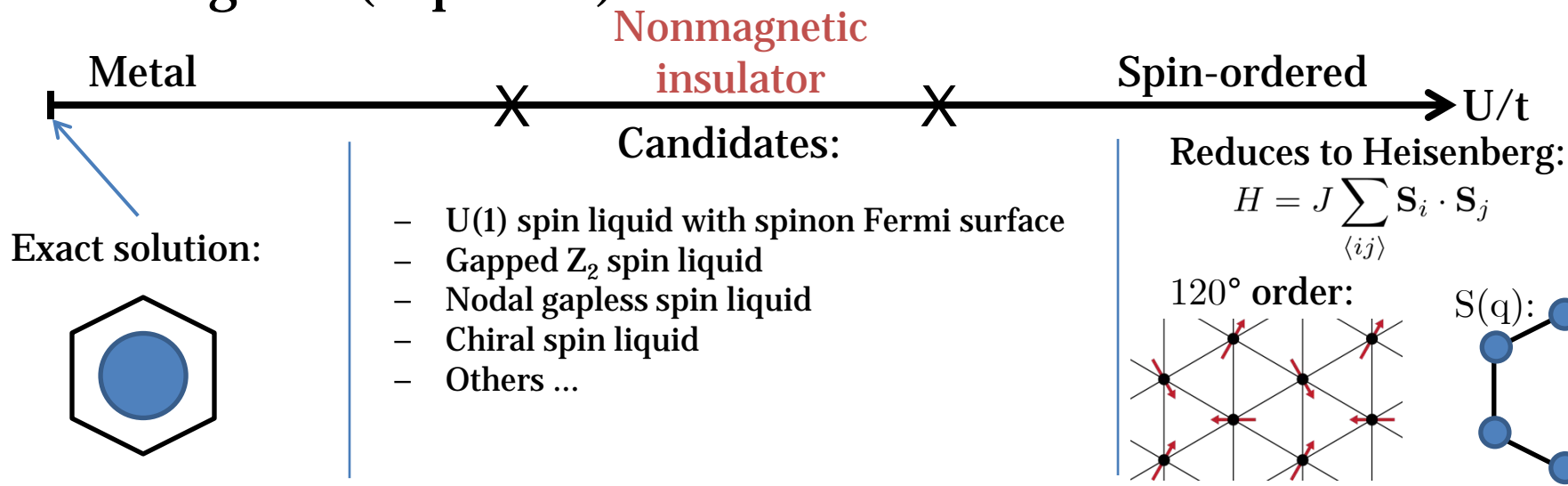
Phase diagram: expected

Hubbard model:

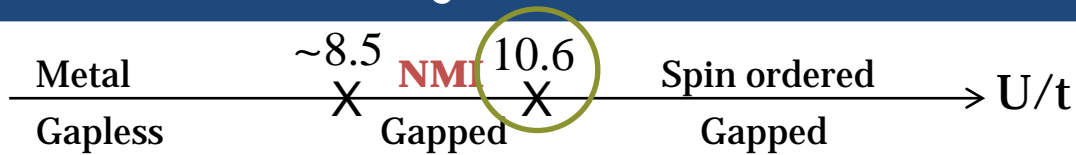
$$H = -t \sum_{\langle i,j \rangle, \sigma} (c_{i\sigma}^\dagger c_{j\sigma} + \text{h.c.}) + U \sum_i n_{i\uparrow} n_{i\downarrow}$$



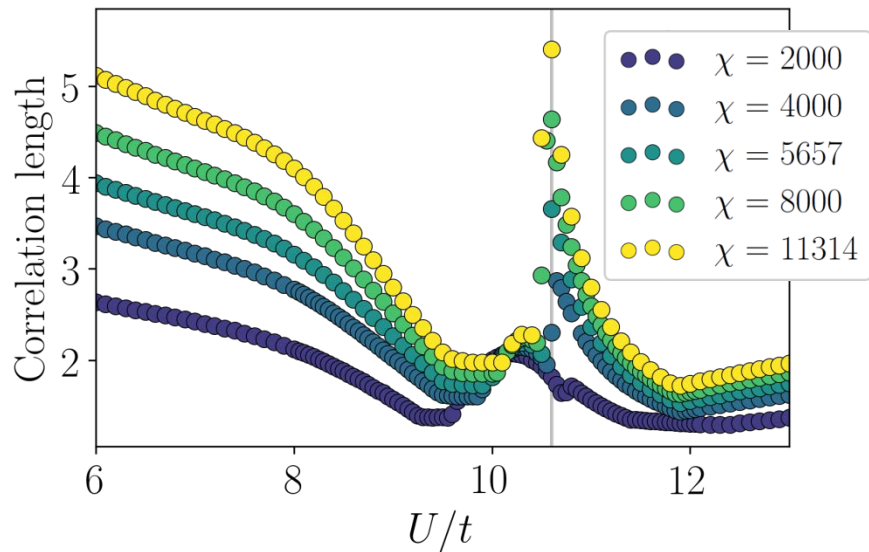
Phase diagram (expected):



Phase diagram: L=4 cylinder



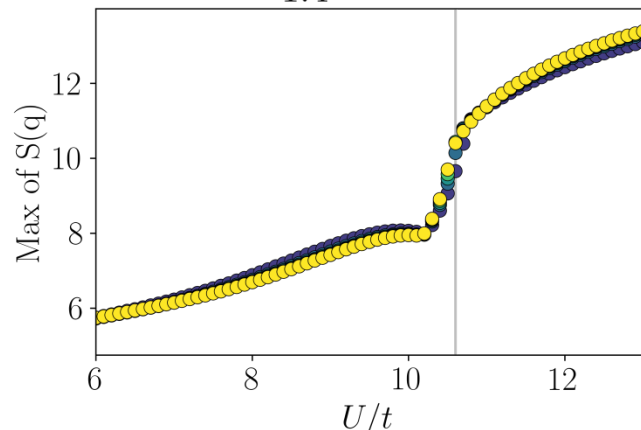
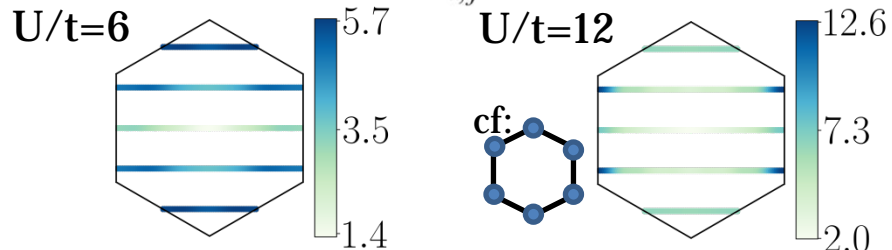
Correlation length



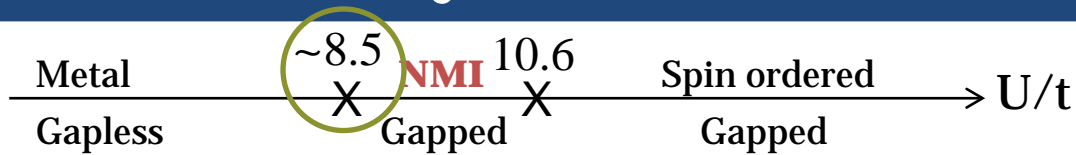
χ : MPS bond dimension

- Controls precision of DMRG

$$\text{Spin order: } S(q) = \frac{1}{N_s} \sum_{i,j} e^{iq \cdot (R_i - R_j)} \langle S_i \cdot S_j \rangle$$



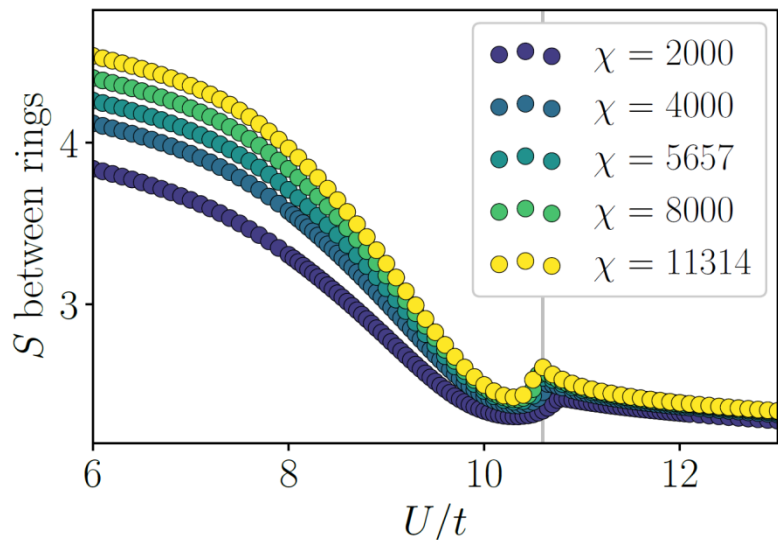
Phase diagram: L=4 cylinder



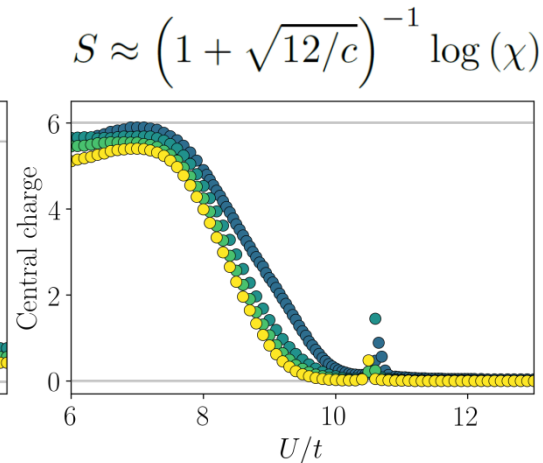
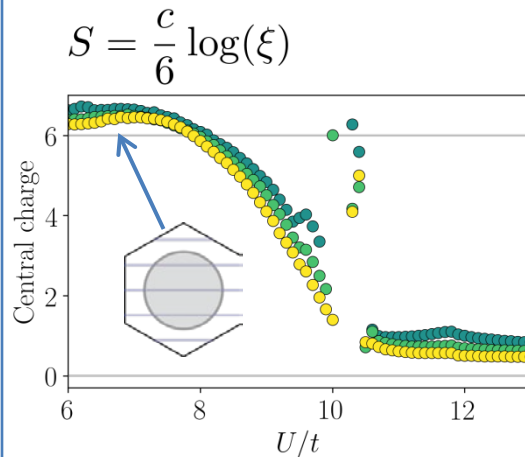
Entanglement:

Diagram illustrating entanglement on a cylinder. A cylinder is shown with a green slice between regions A and B. The entangled state is given by:

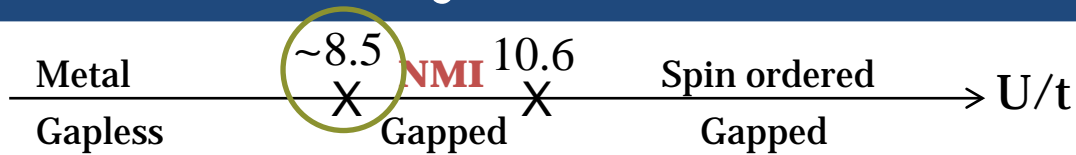
$$|\psi\rangle = \sum_{i=1}^{\infty} \chi \lambda_i |\psi_i^A\rangle |\psi_i^B\rangle \rightarrow S = \sum_{i=1}^{\infty} -\lambda_i^2 \log(\lambda_i^2)$$



Finite entanglement scaling:



Phase diagram: L=4 cylinder

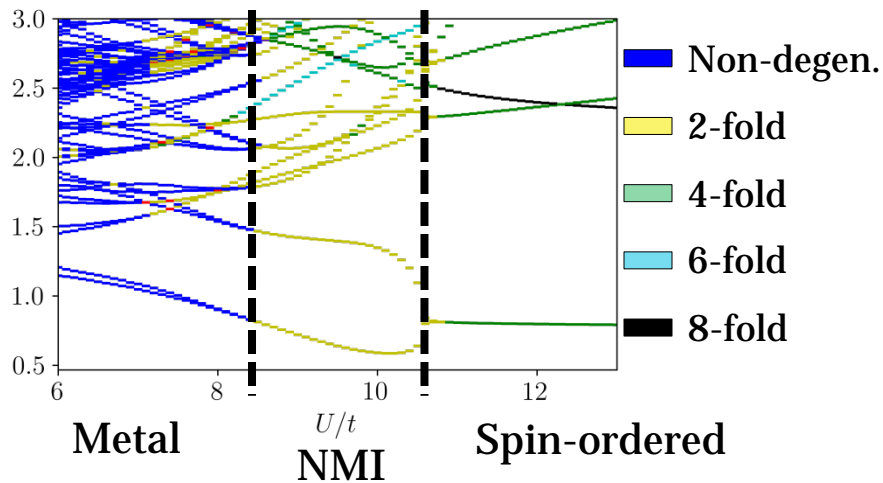


Entanglement:

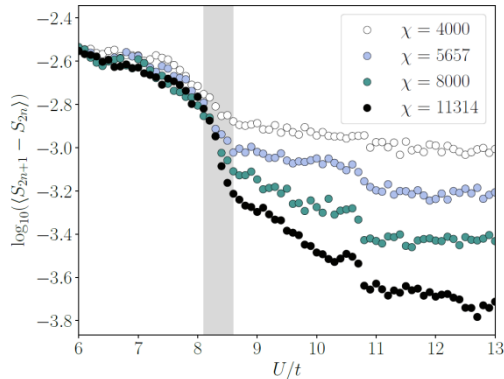
Diagram illustrating the entanglement spectrum. A cylinder is shown with a green slice removed, labeled A and B. The entangled state $|\psi\rangle$ is represented as a sum over eigenstates $|\psi_i^A\rangle$ and $|\psi_i^B\rangle$ with coefficients λ_i . The entanglement entropy S is given by the formula:

$$|\psi\rangle = \sum_{i=1}^{\infty} \lambda_i |\psi_i^A\rangle |\psi_i^B\rangle \quad \rightarrow \quad S = \sum_{i=1}^{\infty} -\lambda_i^2 \log(\lambda_i^2)$$

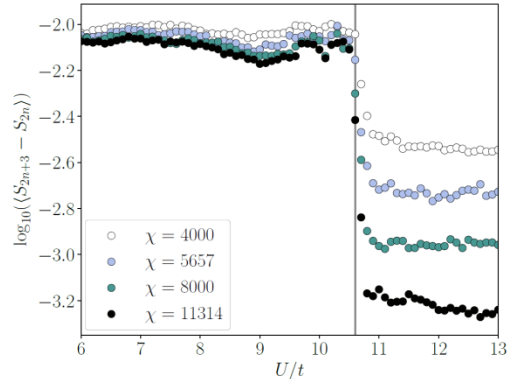
Entanglement spectrum: $\{-\log(\lambda_i)\}$



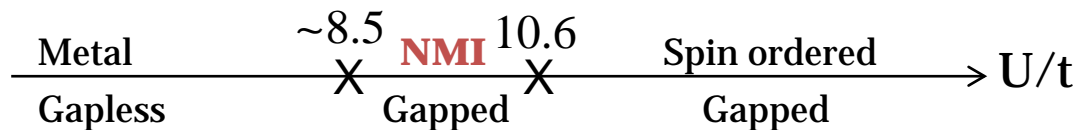
Separation of pairs



Separation of "quads"

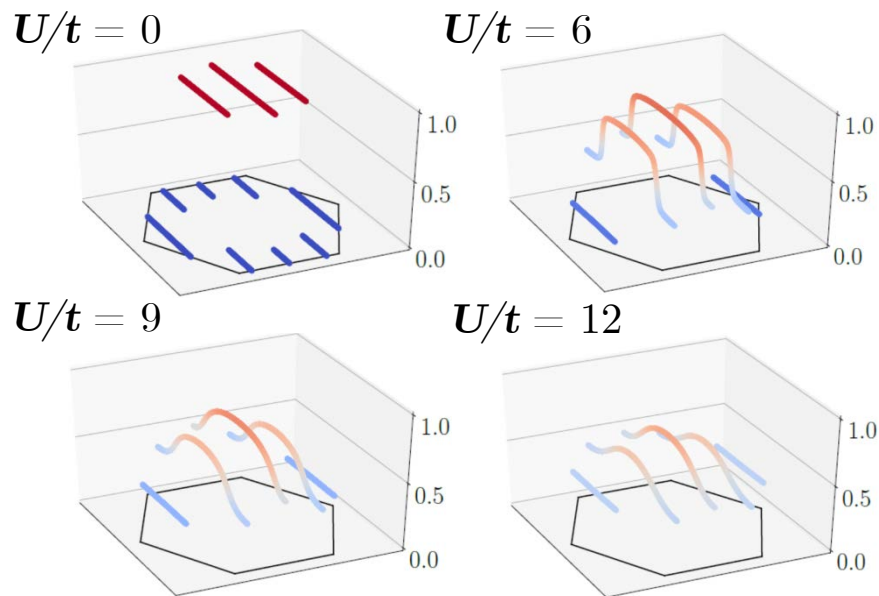


Phase diagram: L=4 cylinder



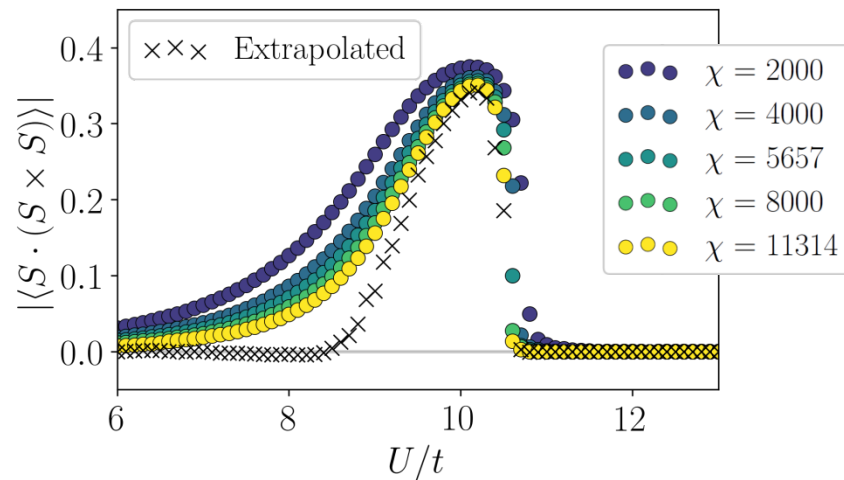
Occupation and Fermi surface:

$$\langle n_{k_x, k_y, \uparrow} \rangle = \sum_{x=-50}^{50} e^{ik_x x} \langle c_{0, k_y, \uparrow}^\dagger c_{x, k_y, \uparrow} \rangle$$

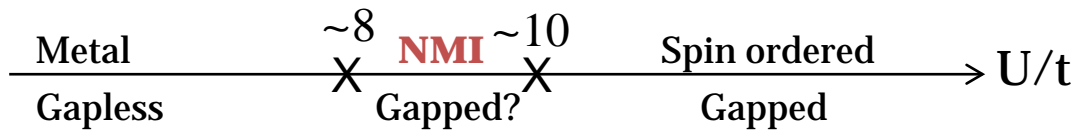


Scalar chiral order parameter:

$$\langle \mathbf{S}_i \cdot (\mathbf{S}_j \times \mathbf{S}_k) \rangle$$



Phase diagram: L=6 cylinder

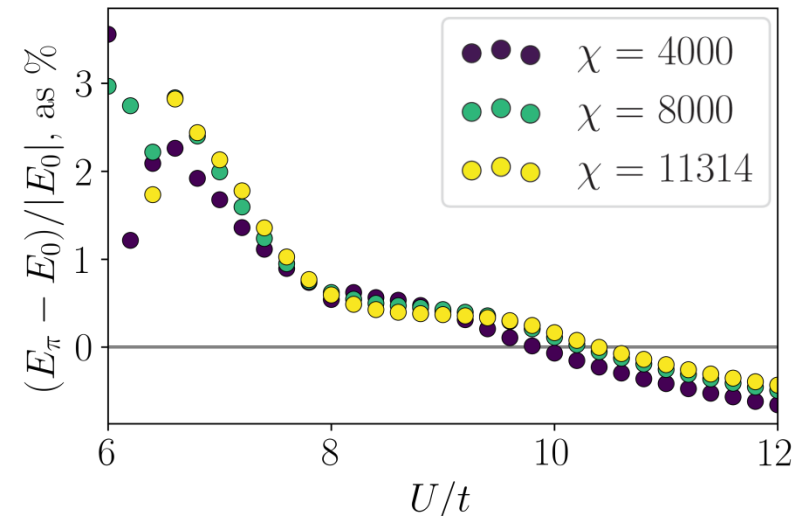
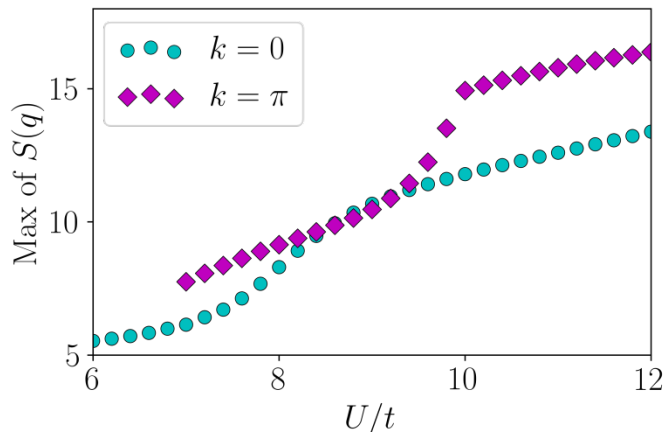
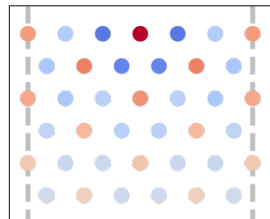
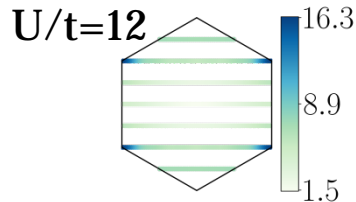
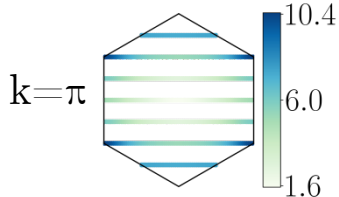
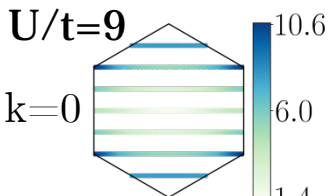
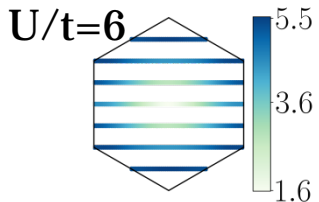


Two low-energy states:



$k=0, k=\pi$
per ring

Spin order:

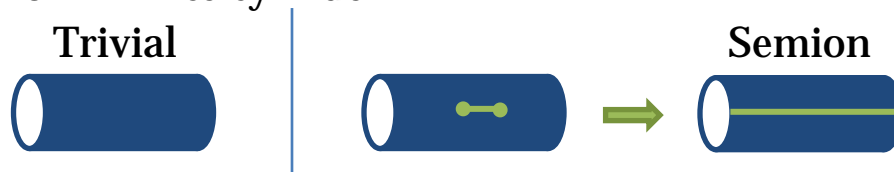


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What is a chiral spin liquid?

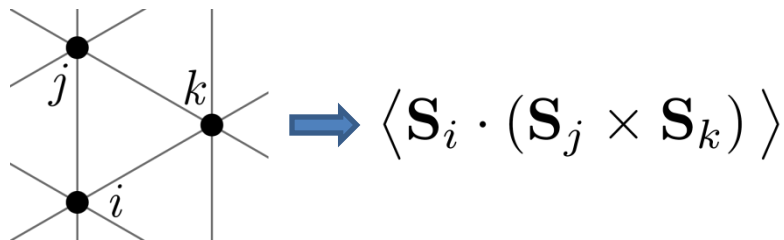
- $\nu=1/2$ fractional quantum Hall effect state for spins
- Signatures:
 - Time-reversal symmetry breaking
 - eg: scalar chiral order parameter
 - 2 chiralities \rightarrow 2x ground state degeneracy
 - Topological ground state degeneracy and fractionalized quasiparticles (semions)
 - 2x on infinite cylinder



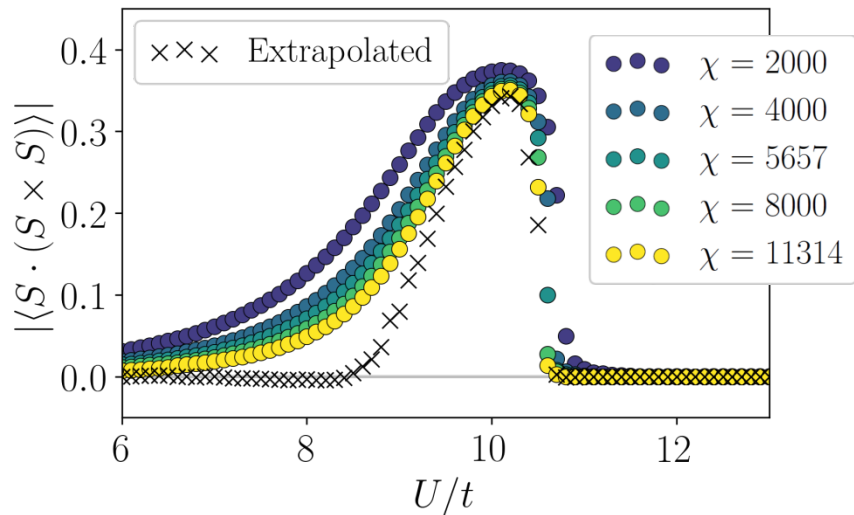
- Chiral edge modes in 2D \rightarrow entanglement spectrum on cylinder
 - Characteristic level counting vs momentum: 1, 1, 2, 3, 5, ...
 - 2x degeneracy for semion sector
- Quantized spin Hall effect: 2π flux insertion \rightarrow spin $1/2$ pumping

Identification as a CSL

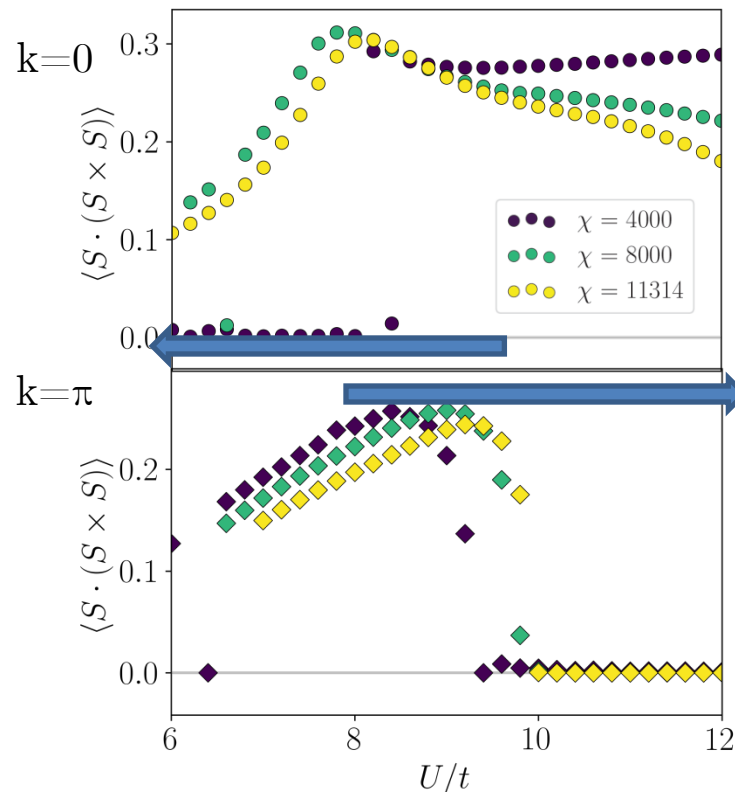
Chiral order parameter:



L=4



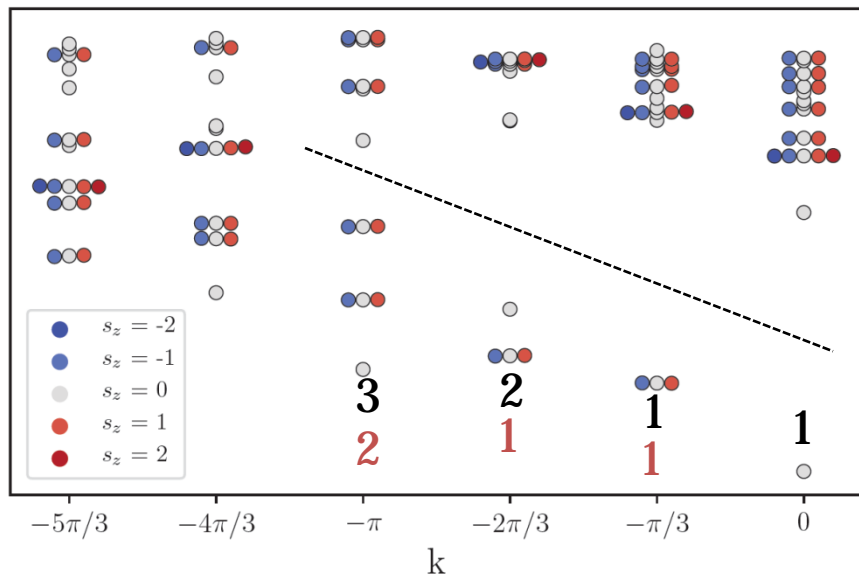
L=6



Identification as a CSL

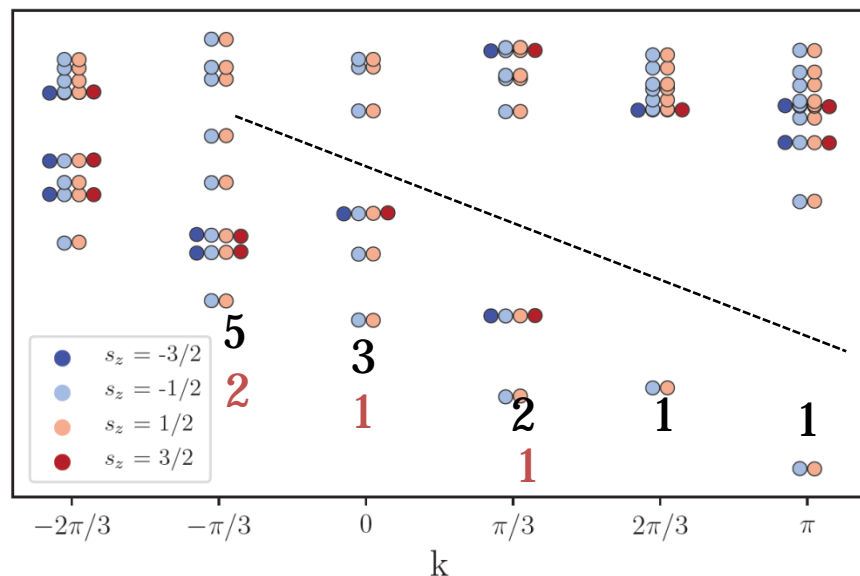
Spin- and momentum-resolved entanglement spectrum, $L=6$, $U/t = 9$:

$k=0$



Trivial sector

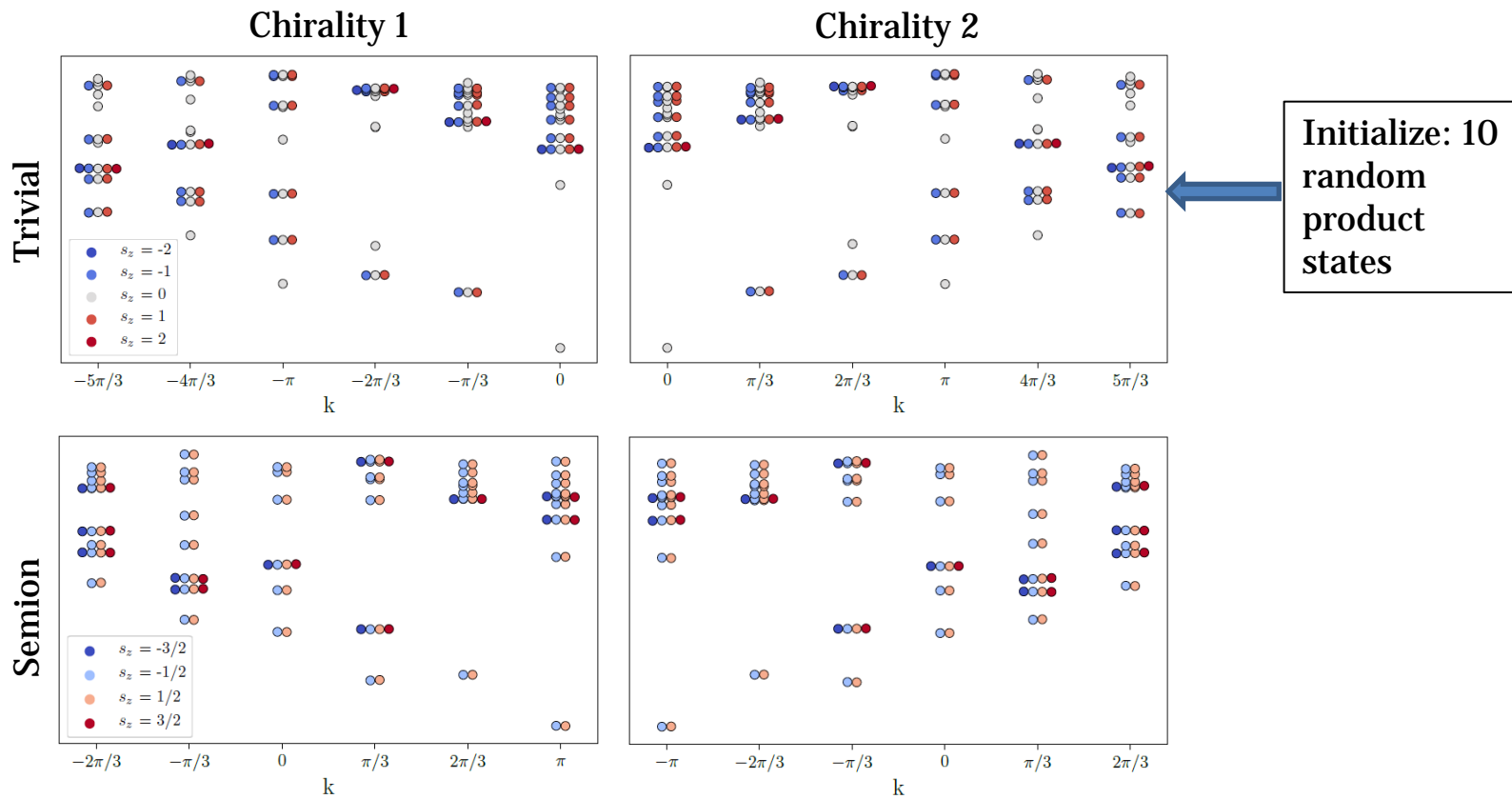
$k=\pi$



Semion sector

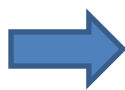
Identification as a CSL

Ground state degeneracy, $L=6$, $U/t = 9$:



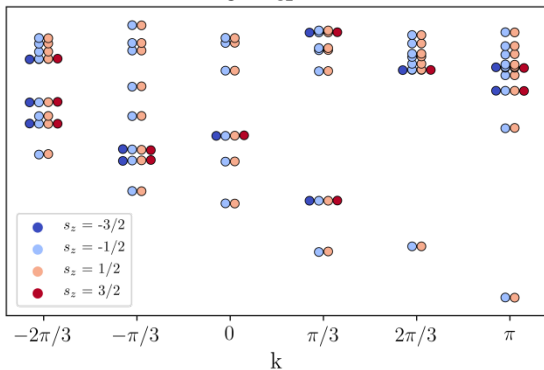
Identification as a CSL

Flux insertion and spin Hall effect:

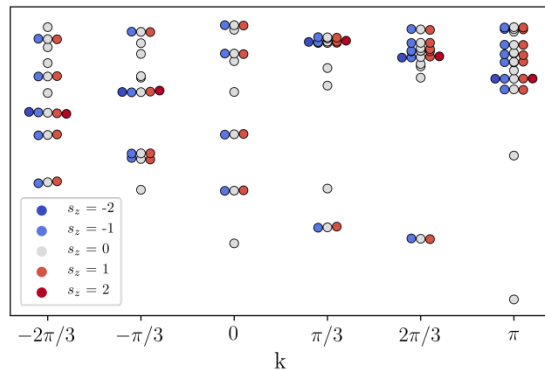


$$H \rightarrow -t \sum_{\langle ij \rangle} \left(e^{i(i_y - j_y)\sigma\theta/2} c_{i\sigma}^\dagger c_{j\sigma} + \text{H.c.} \right) + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

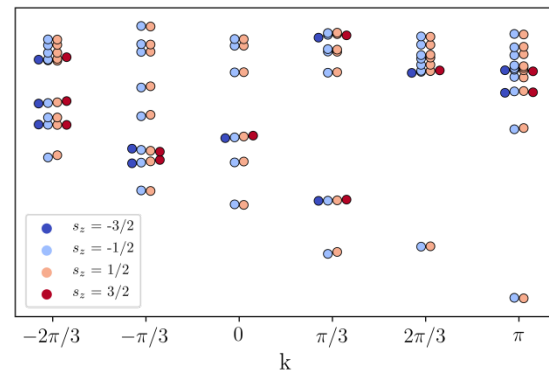
0 flux



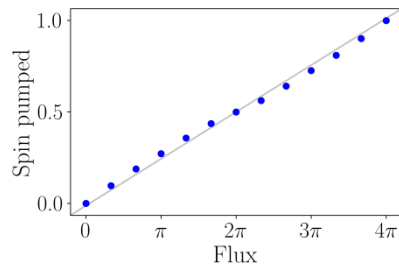
2π flux



4π flux



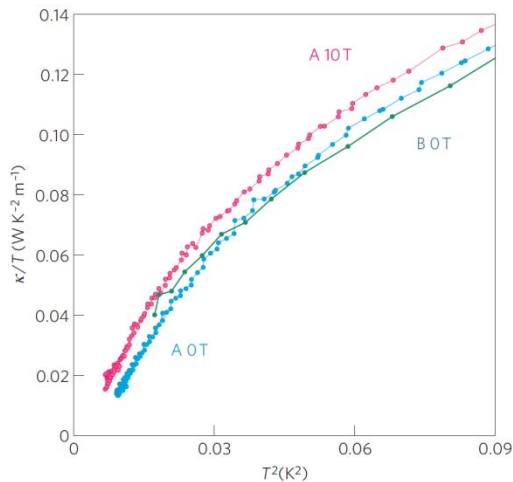
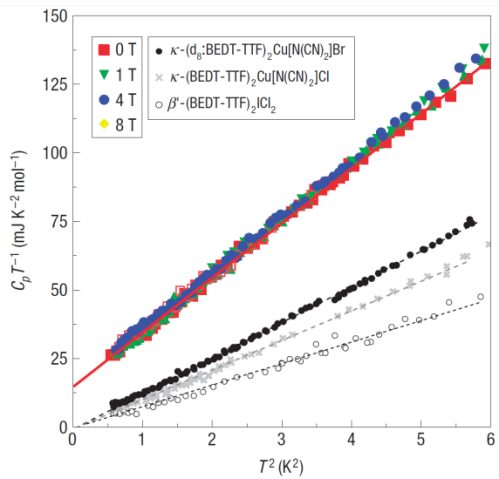
Spin Hall effect:



Outline

1. Introduction
2. Calculation methods
3. Phase diagram
4. Chiral spin liquid phase
- 5. Implications/comparisons and summary**
- 6. Future directions**

Comparison with experiments

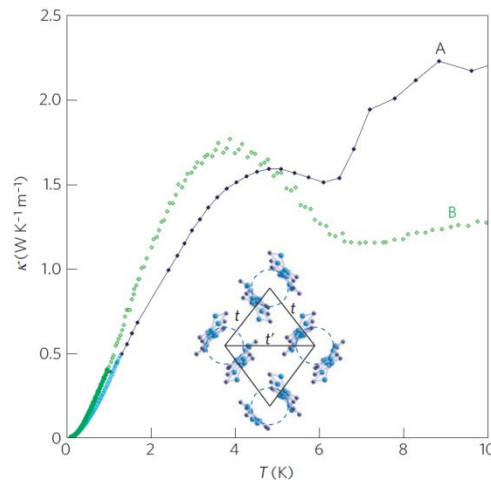


Gapless heat capacity

Gapped conductivity

Gapless chiral edge modes?

Peak in thermal conductivity

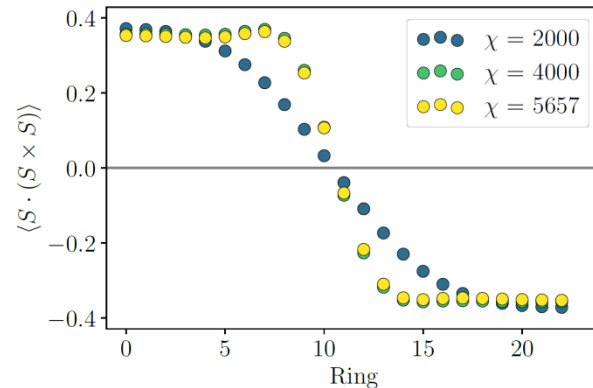
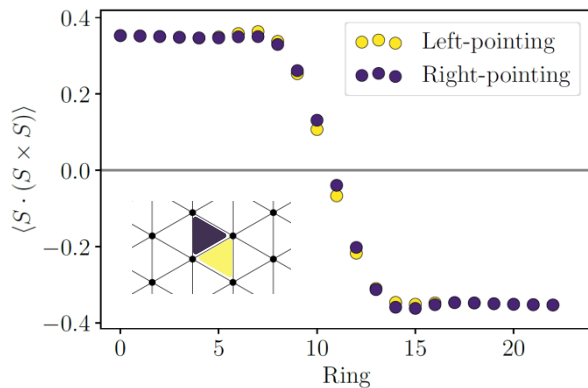
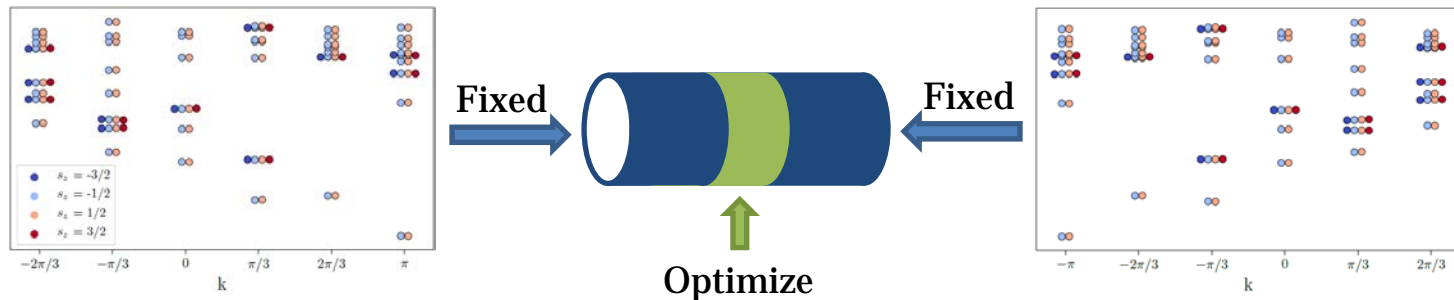


Maybe a finite T phase transition?

Ising-like ordering transition for two chiralities?

Look at chiral domain wall tension

Chiral domain wall



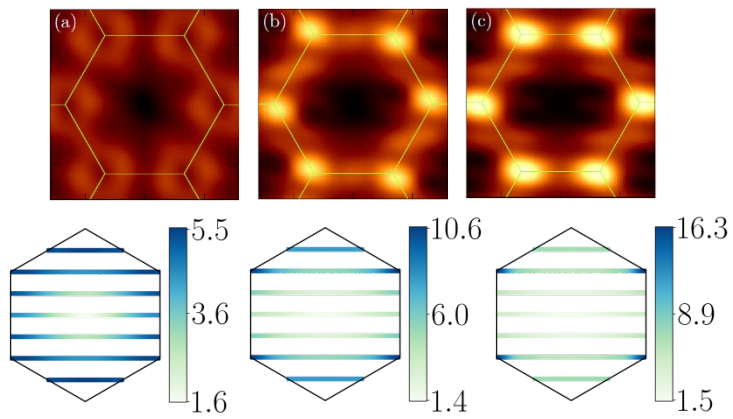
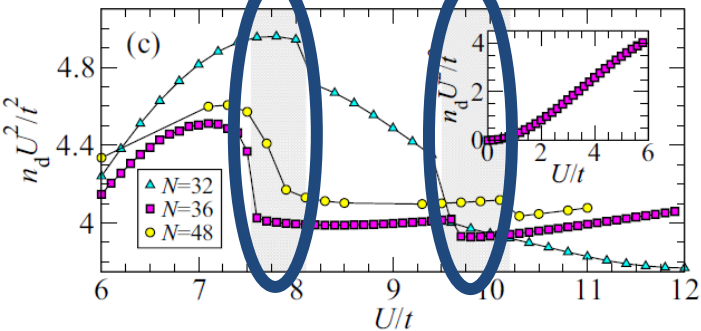
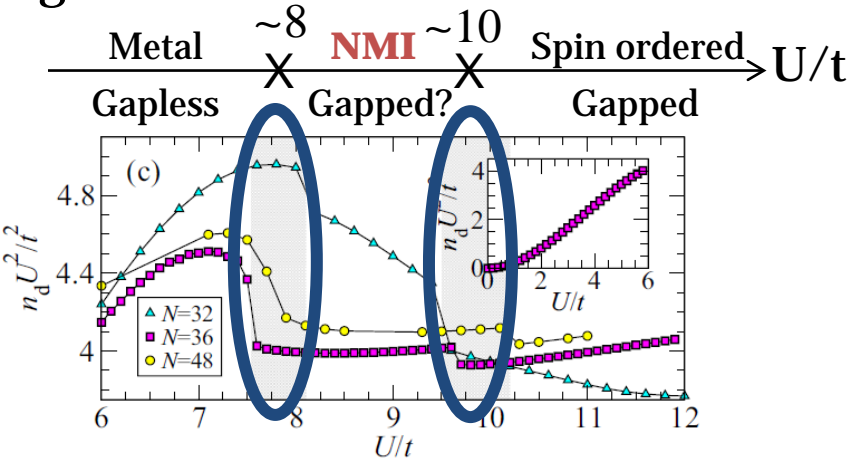
Domain wall tension:
 $0.0065 t/a \approx 4 K$



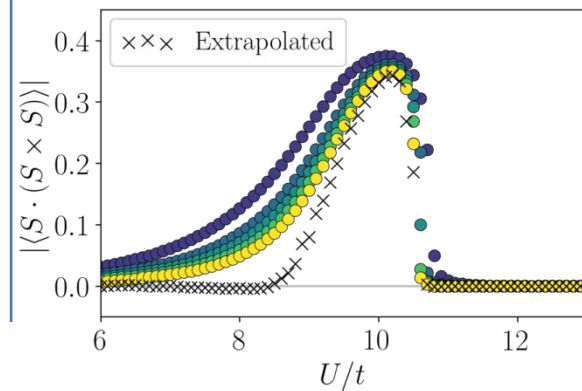
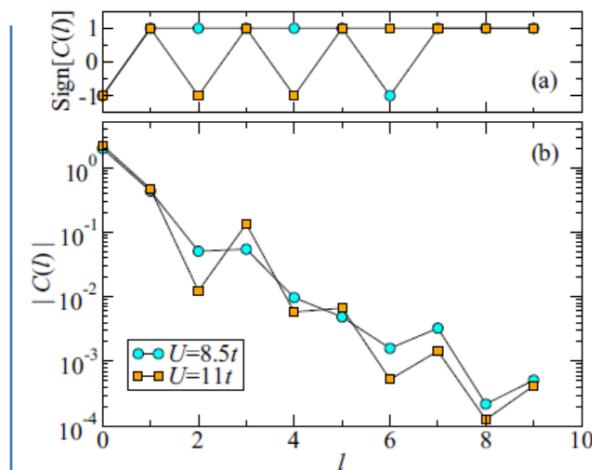
Same order of magnitude:
 may explain transition

Comparison with RIKEN group results

Agreement:



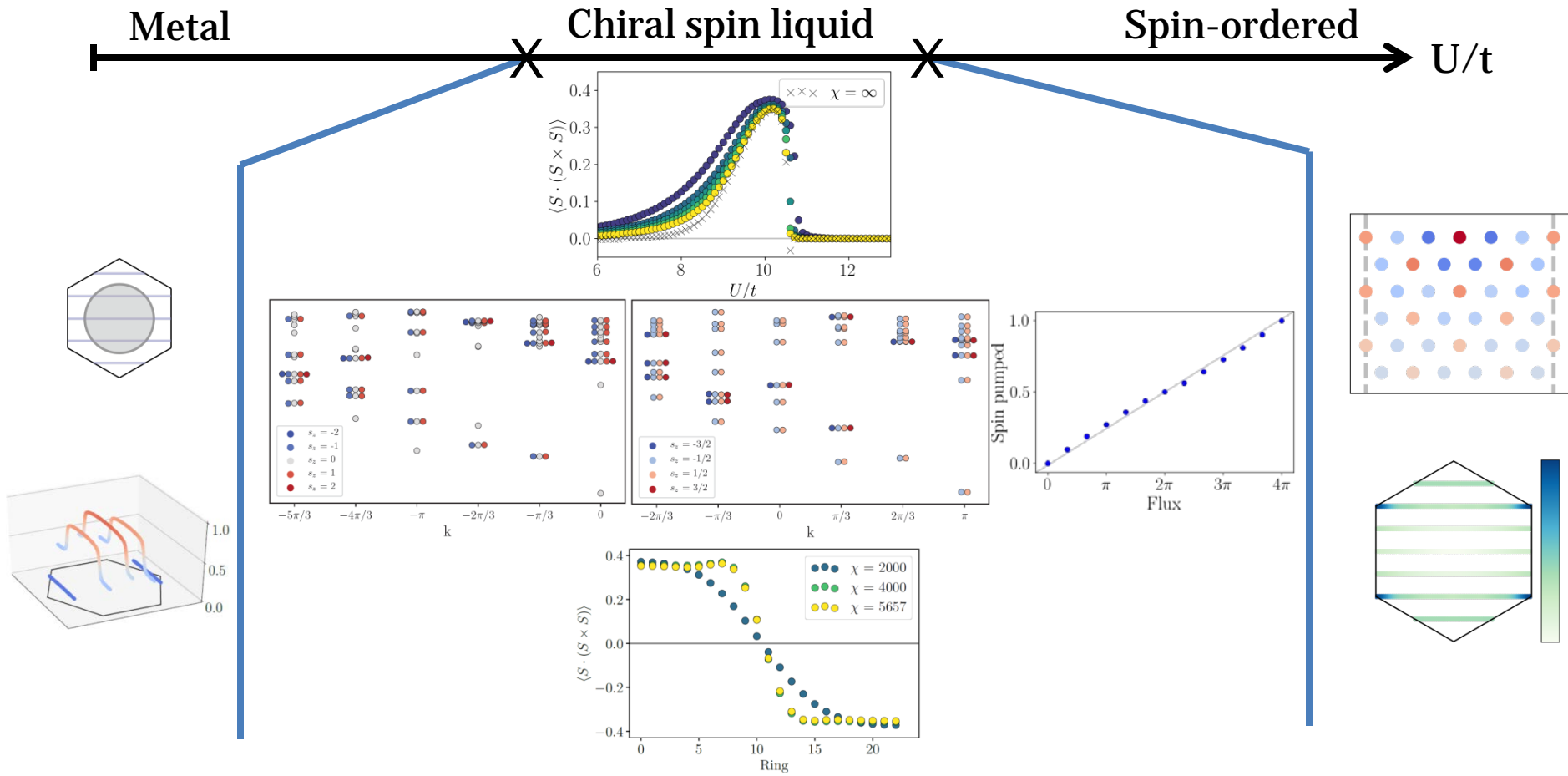
Disagreement:



Why?

- Different bcs?
- k-space?
- Finite vs infinite?

Summary

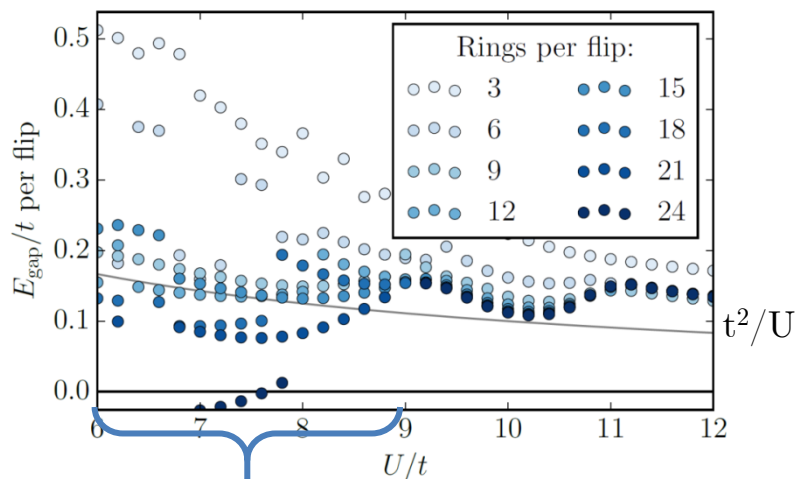
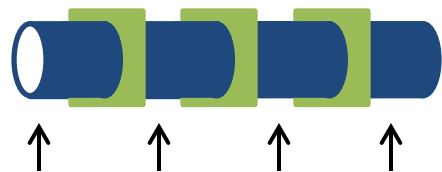


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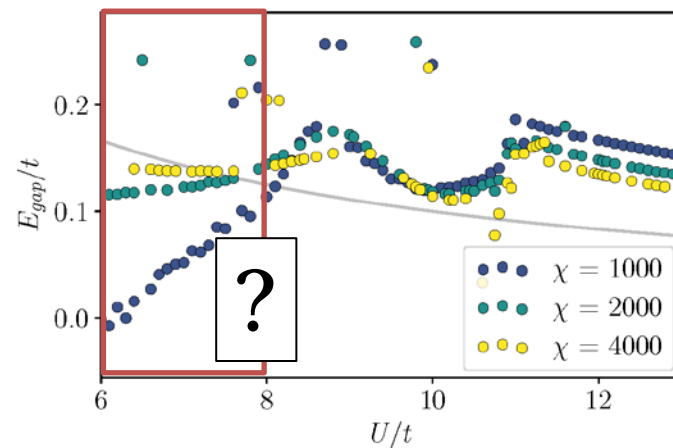
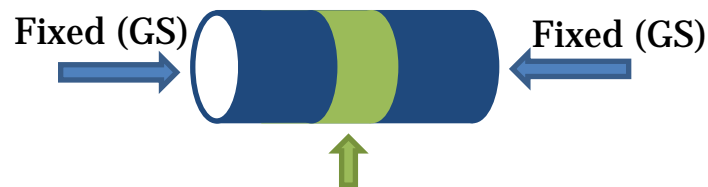
Spin gap

Version 1: Flip 1 spin per N rings:



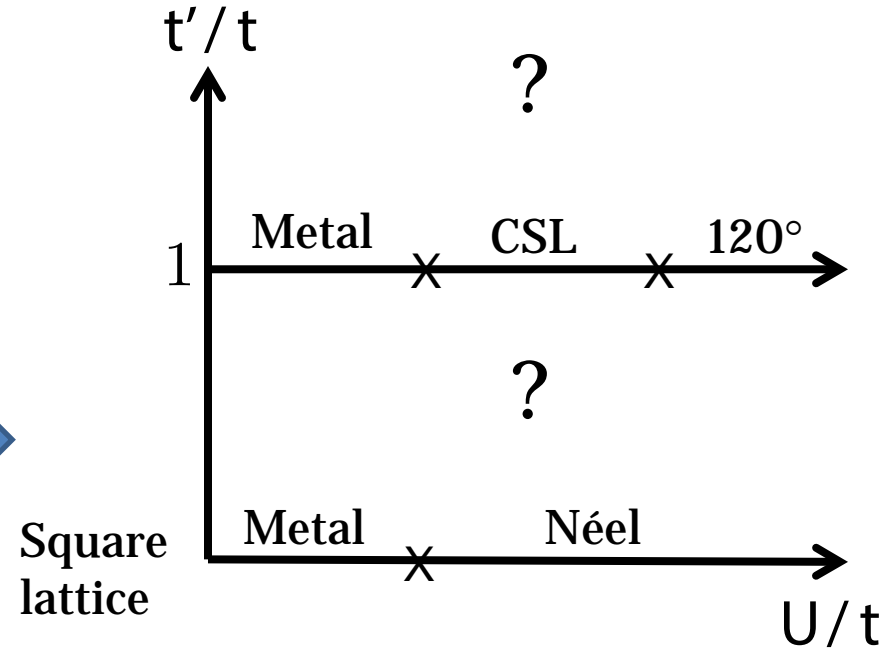
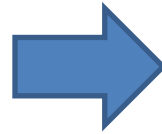
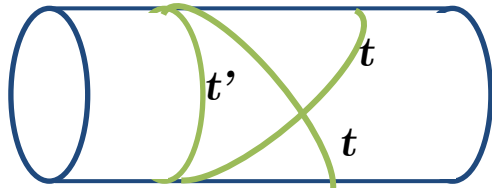
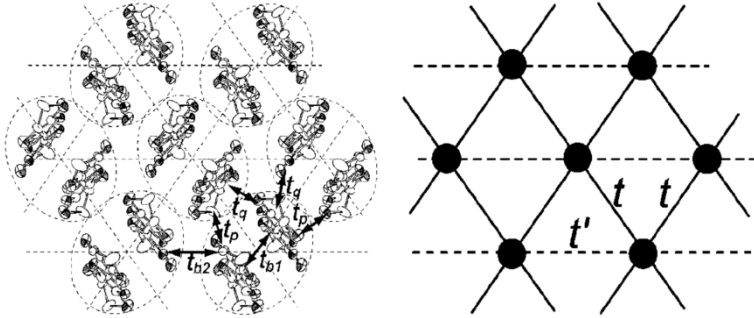
???: Take it with a grain of salt!

Version 2: Fixed edges



Anisotropy

Real material has some anisotropy:



- Interesting phase transitions to observe!
- Insight for isotropic case phase transitions?

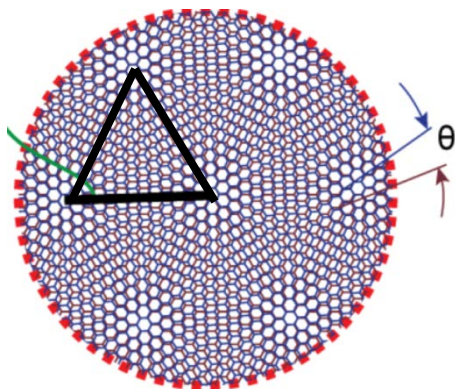
Mixed-space DMRG

For triangular lattice:

- Extra quantum number \rightarrow computational efficiency
- Find ground state in specific momentum sector

Can be even more useful!

Consider system with Moiré pattern
eg. twisted bilayer graphene



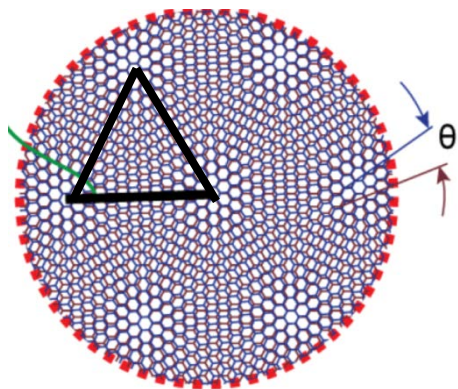
Cao et al., Nature 2018

Extremely large
unit cell

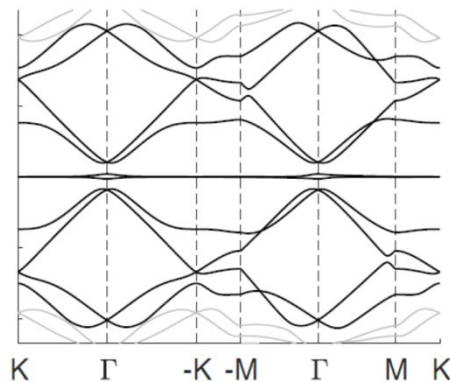


~~DMRG~~

Mixed-space DMRG



Cao et al., Nature 2018



Po et al., 1808.02482

Only need flat bands

Use Wannier states (k_y, x)

$$\{t_{k_1 k_2, x_1 x_2}\}$$
$$\{V_{k_1 k_2 k_3 k_4, x_1 x_2 x_3 x_4}\}$$

Automatically generate and compress MPO

DMRG

Acknowledgements

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(UC Berkeley)



Joel E. Moore
(UC Berkeley)

TenPy DMRG code:

- Michael Zaletel
- Frank Pollmann (Munich)
- Roger Mong (Pittsburgh)

Computing resources:

- Lawrence Berkeley National Laboratory



Summary

- Three phases of Hubbard model on triangular lattice
 - Metal, nonmagnetic insulator (NMI), magnetically ordered
- NMI phase is a chiral spin liquid!
 - Chiral order parameter → spontaneous breaking of time-reversal symmetry
 - Two topologically degenerate ground states: trivial, semion sectors
 - Spin Hall effect: 2π flux insertion pumps spin $\frac{1}{2}$
 - May explain features observed in experiments

For more see arXiv: 1808.00463

Thanks for your attention!